

Electrical characterization of the backside interface on BSI global shutter pixels with Tungsten-shield test structures on CDTI process.

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Abstract-A new methodology is presented using well known electrical characterization techniques on dedicated single devices in order to investigate backside interface contribution to the measured pixel dark current in BSI CMOS image sensors technologies. Density of interface states and charges within the oxide extractions are achieved. The results show that in our case, the density of state is not directly the source of dark current excursions. The quality of the passivation of the backside interface appears to be the root cause. Thanks to the presented new test structures, it has been demonstrated that backside interface contribution to dark current can be investigated separately from other contributions.

I.INTRODUCTION

Backside illuminated (BSI) technologies have evolved to incorporate new backside interface passivation techniques (i.e. high-K negatively charged dielectric stack to easily accumulate holes and therefore passivate the interface [1]) and tungsten (W) layer for shielding requirement of BSI Global Shutter pixels. Interface states and/or the quality of the passivation can be key for dark current. Therefore it is crucial to develop this interface characterization. Figure 1 shows the median distributions of dark current (I_{dark}) measured on a pixel array at wafer level for three wafers processed with a BSI CDTI technology, the dielectric stack being a Ta₂O₅/Al₂O₃ with process condition variants. Atypical behaviors can be seen on wafers 2 and 3. For both wafers, there is a higher and non-uniform dark current. In the following work, the different measurements have been performed on the same dies as for the I_{dark} measurement (Figure 2). To study the potential correlation between I_{dark} , density of interface states (D_{it}) and quality of passivation (negative charges within the oxide N_{eff}), two test structures have been developed in a BSI global shutter pixel technology (with CDTI [2]) using the W-shield as a gate: a backside MOS capacitor (Figure 3) ($L = 500 \mu m$, $W = 31.6 \mu m$) and a backside W-shield gate pseudo-MOS transistor (Figure 4) ($L = 10 \mu m$, $W = 27 \mu m$). In the last structure, CDTI are biased in a way that there is always an inversion layer at their interface in order to connect surface source and drain to the bottom interface of interest (channel of the W-shield gate transistor). Figure 5 shows the conventional operation of the transistor test structure. In the next sections, first D_{it} is investigated with a tentative correlation with I_{dark} . Then the passivation quality is

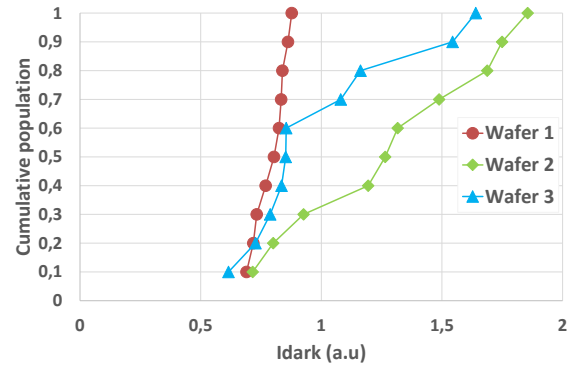


Figure 1 : Cumulative population of the pixel dark current measured on testchips on several dies of 3 wafers showing distinct behaviors depending on the wafer.

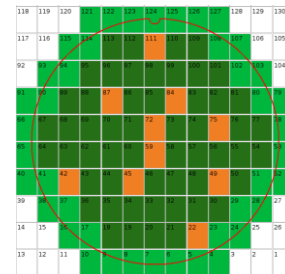


Figure 2 : Sampling map of the measured dies (in orange) for all the measurements.

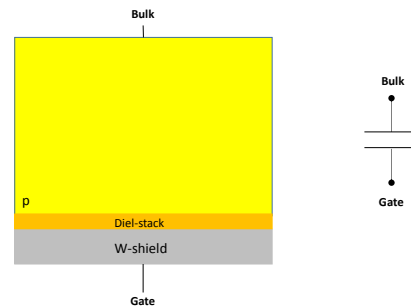


Figure 3 : Illustration of the capacitive test structure used to extract N_{eff} .

studied. Finally, a charging phenomenon present on both structures is presented.

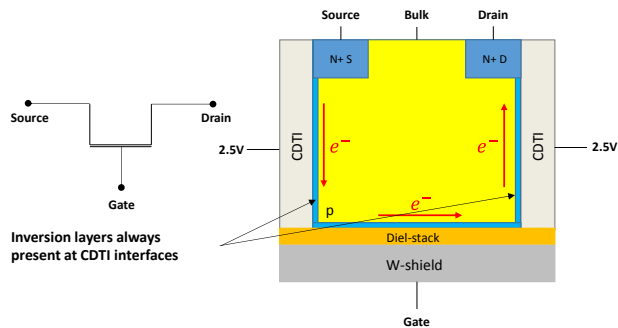


Figure 4: Illustration of the transistor test structure used to extract D_{it} (an isolation is present between the source and drain to prevent punchthrough).

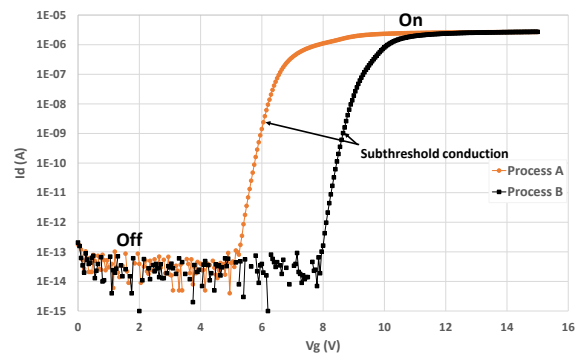


Figure 5 : $I_d(V_g)$ curves showing the conventional operation of the atypical transistor with clear ON/OFF behaviors.

II. INTERFACE STATES CHARACTERIZATION

The dark current generated by an interface may be dominated by interface states if their density is high. In order to extract the interface states density (D_{it}) for each die, the charge pumping method [3] [4] [5] is applied on the W-shield transistor test structure (Figure 4). The method is illustrated by Figure 6. The CDTI are biased at 2.5V in order to have an inversion layer at their interface. The drain and the source are grounded. A square pulse with the following characteristics is applied to the W-shield gate: $V_{gpulsebase}$ is swept from 2V to 12V, $\Delta V_{gpulse} = 6V$, minimum frequency of the signal $f = 50$ kHz and rise and fall time $T_{r,f} = 1\mu s$. The pumped current is measured through the bulk contact. The choice of the minimum frequency was made in order to avoid the measurement of traps in the High-K dielectric volume [6]. Using the method described in [7], no geometric effect was detected in the measurements. The maximum measured pumped current is proportional to the mean D_{it} (cm^{-2}):

$$I_{pump_{max}} = D_{it}qSf \quad (1)$$

Here, $I_{pump_{max}}$ represents the maximum pumped current, q is the elementary charge, S the gate area. Figure 7 shows the maximum pumped current as a function of the

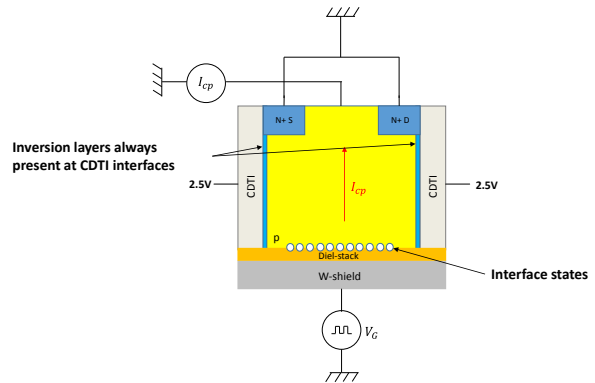


Figure 6: Charge pumping method applied to W-shield transistor test structure.

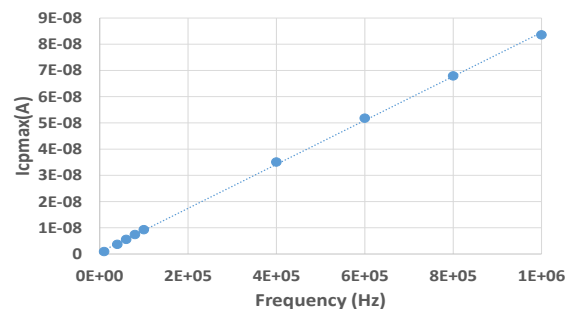


Figure 7: $I_{cpmax}(f)$ curves with the transistor test structure.

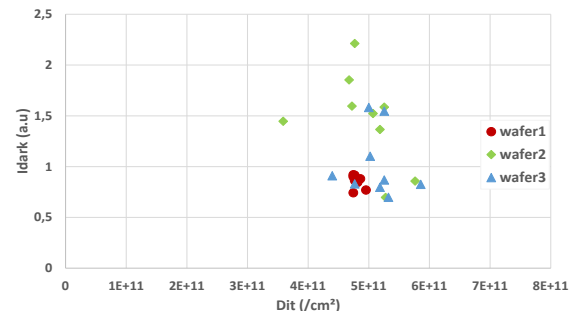


Figure 8 : $I_{dark}(D_{it})$ scatter plot. All the wafers show a same level of D_{it} .

frequency. It can be seen that the maximum measured pumped current is linear with f which is in agreement with equation (1). The slope extraction gives D_{it} . Measurements are performed on the 3 wafers using the sampling map of Figure 2. Figure 8 presents the I_{dark} vs D_{it} scatter plot. It does not show any clear tendency. D_{it} looks not to be correlated to I_{dark} in this case. So the measured dark current differences between the 3 wafers do not seem to be explained by a difference of D_{it} , it can be seen that all the wafers have an equivalent level of interface states ($\sim 5 \times 10^{11}/cm^2$).

III. QUALITY OF THE PASSIVATION CHARACTERIZATION

As interface states do not seem to be the response for I_{dark} excursions, the quality of the passivation

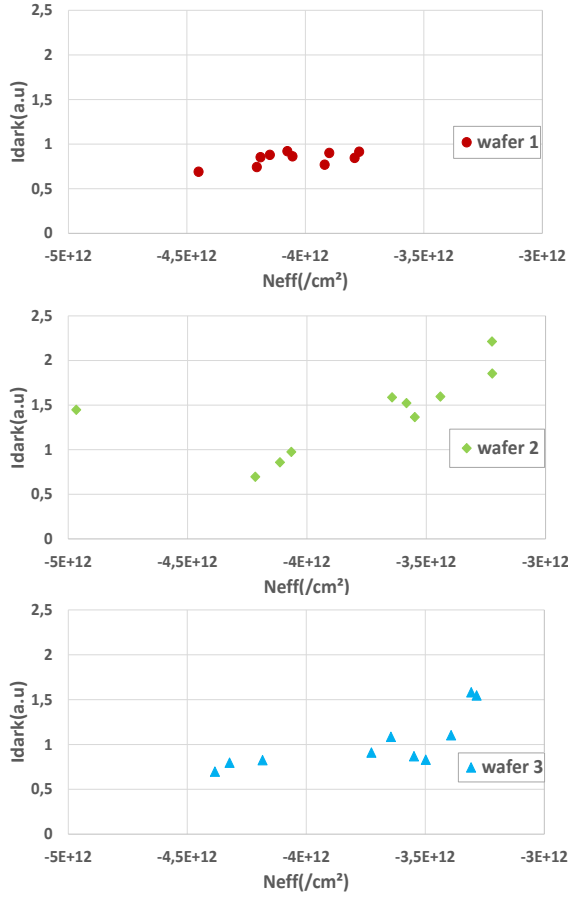


Figure 9 : $I_{dark}(N_{eff})$ scatter plot for the 3 wafers.

might be another root cause of the I_{dark} response seen in Figure 1. Indeed, with very similar Dit on the measured wafers, the field effect passivation, that is, having an electric field at the interface to accumulate charges in order to passivate the interface, can have a key contribution. The electric field is induced by negative charges within the oxide (N_{eff}). On the following analysis the Maserjian's function is used on C-V measurements to estimate the effective charge density within the oxide on the capacitive structure (Figure 3) [8] [9]:

$$Y(V_g) = \left(\frac{1}{C}\right)^3 \frac{dC}{dV_g} \quad (2)$$

Thanks to this function, the substrate doping concentration (N_a) and the flat band voltage (V_{FB}) can be extracted:

$$Y_{min} = -\frac{1}{q\epsilon_{Si}N_a} \quad (3)$$

$$Y(V_{FB}) = -\frac{Y_{min}}{3} \quad (4)$$

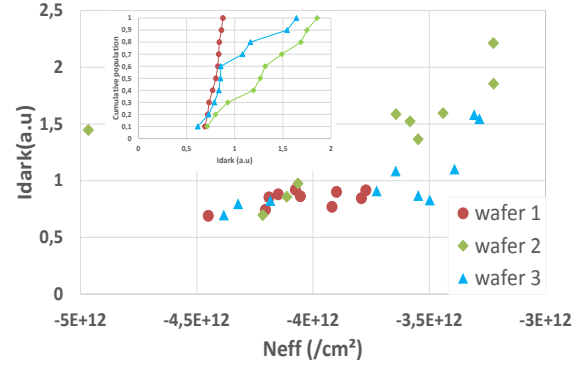


Figure 10 : $I_{dark}(N_{eff})$ scatter plots for wafers 1, 2 and 3. Contrary to the Dit study, a correlation is visible between N_{eff} and the measured dark current.

Where Y_{min} is the minimum reached by the Maserjian's function and ϵ_{Si} the silicon permittivity. Then it is possible to calculate N_{eff} :

$$N_{EFF} = \frac{-(V_{FB} - W_{MS})C_{ox}}{q} \quad (5)$$

Where N_{EFF} is the effective number of charges within the oxide, W_{MS} the metal semiconductor work function and C_{ox} the oxide capacitance. For the measurements, the following characteristics are applied: V_g is swept from 2V to 12V, the signal frequency applied is $f = 50$ kHz, and the modulation of the signal is 0.02V. Figure 9 presents the scatter plots of I_{dark} vs extracted N_{eff} for the 3 wafers. On wafer 1, it is hard to extract an eventual correlation probably because the I_{dark} distribution is very narrow (Figure 1). However, on wafers 2 and 3, a clear correlation can be identified between N_{eff} and I_{dark} . In order to observe if there is a more general tendency, a global scatter plot with the 3 wafers together is shown in Figure 10. On this figure, two parts of the scatter plot can be distinguished: a first one where the dark current increase according to N_{eff} is weak (almost a plateau) and a second one where the dark current increase according to the N_{eff} decrease (in absolute value) is clearly visible. According to [1], the plateau is explained by the fact that when a certain number of charges within the oxide is reached, the interface is fully passivated and the I_{dark} from the backside interface becomes negligible and the measured dark current comes from another part of the pixel (DTI, frontside interface, silicon volume etc.). On the contrary, when N_{eff} decrease (in absolute value) as on some wafer 2 and 3 dies, the passivation is not efficient enough and a significant I_{dark} appears coming from the BS interface.

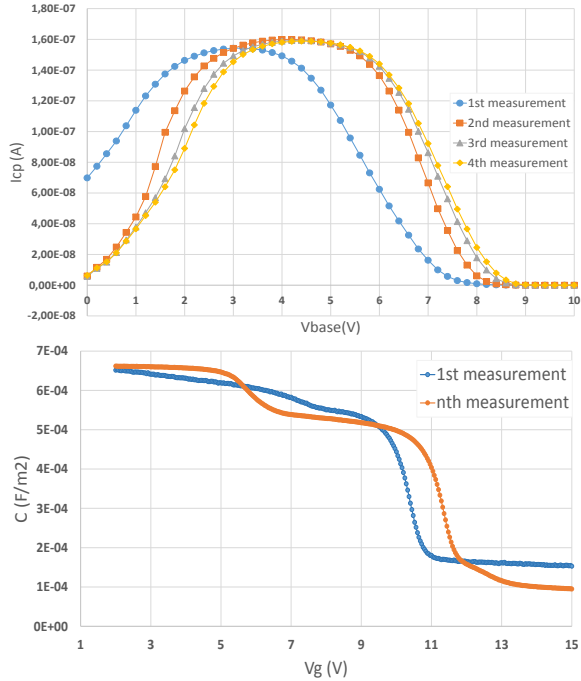


Figure 11 : Charging effect measured on both MOS and W-shield transistor structures.

IV. CHARGING EFFECT ON BOTH STRUCTURE

After a first measurement up to relatively high voltage on both structure, a charging effect occurs when the structures are measured a second time and again. This effect can be seen on Figure 11. In both cases, the curve shift seems to be related to a V_{FB}/V_t shift. In the case of the charge pumping, between the first measurement and the other ones, the maximum pumped current slightly increases which is the result of an increase of D_{it} according to equation (1). In the case of the C-V measurement, it can be seen that a more pronounced bump is present after a certain number of measurements. A bump in C-V measurements can be the result of D_{it} centered within the gap [10]. Concerning the curves shift, one possible explanation is that during the measurements V_g reaches high voltage (here 12V) that can induce a negative charge injection which results in higher V_{FB}/V_t voltage. A hysteresis effect may be seen (and may be correlated with the charging effect) on both structures (Figure 12). In this case, V_g is swept from 0V to 12V and then from 12V to 0V and finally from 0V to 12V again. An effect of charging/discharging can be seen. To be more precise, more investigations would be needed and at least cautions have to be taken. The different extractions presented earlier in this paper were made during the first measurement when structures are not already stressed. After waiting a certain time, the structures come back at their initial state. It should

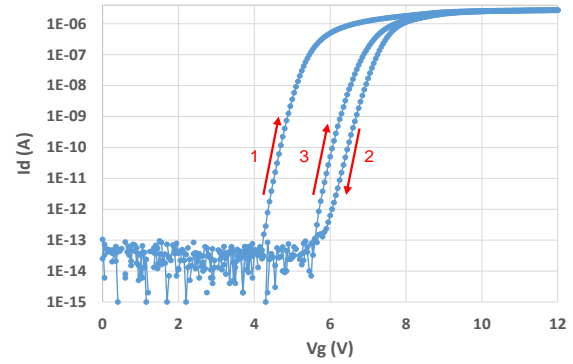


Figure 12: Hysteresis effect on the transistor structure, the same effect can be seen on the capacitance structure.

be noted that for the product utilization, such conditions are not applied and such effect should not happened.

V. CONCLUSION

With these structures, it is possible to electrically characterize the backside interface of a BSI technology using tungsten shield. By means of two known characterization methods, it is possible to extract D_{it} and N_{eff} (which are the two important parameters for dark current). So it is possible to determine if the dark current mainly comes from the backside interface, and to discriminate the origin of the BS dark current. A drawback of this method is the presence of a charging effect which have to be studied more precisely. In addition to these I_{dark} contribution studies, these dedicated devices with associated characterizations can be helpful for process monitoring, TCAD calibration and reliability works.

References:

- [1] G. Dingemans and W. Kessels, "Status and prospects of Al₂O₃-based surface passivation schemes for silicon solar cells," *Journal of Vacuum Science and Technology*, vol. 30, no. 4, pp. 040802-1/27, 2012.
- [2] A. Tournier, F. Roy, Y. Cazaux, F. Lalanne, P. Malinge, M. McDonald, G. Monnot and N. Roux, "A HDR 98dB 3.2 μ m charge domain global shutter CMOS image sensor," in *IEDM*, San Francisco, 2018.
- [3] D. K. Schroder, *Semiconductor material and device characterization*, 3rd ed., Hoboken, New Jersey: John Wiley & Sons, 2006.
- [4] A. B. M. Elliot, "The use of charge pumping currents to measure surface state densities in MOS transistors," *Solid-State Electronics*, vol. 19, pp. 241-247, 1976.
- [5] G. Groeseneken, H. E. Maes, N. Beltran and R. F. D. Keersmaecker, "A Reliable approach to Charge-pumping measurements in MOS transistors," *IEEE Transactions on electron devices*, Vols. ED-31, pp. 42-53, 1984.
- [6] R. Degraeve, A. Kerber, P. Roussel, E. Cartier, T. Kauerauf, I. Pantisano and G. Groeseneken, "Effect of bulk trap density on HFO₂ reliability and yield," in *International Electron Devices Meeting*, 2003.
- [7] L. Lu, Mingxiang and M. Wong, "Geometric Effect Elimination and Reliable Trap State Density Extraction in Charge Pumping of Polysilicon Thin-Film Transistors," *IEEE Electron Device Letters*, vol. 30, no. 5, pp. 517-519, 2009.
- [8] J. Maserjian, G. Petersson and C. Svensson, "Saturation capacitance of thin oxide MOS structures and the effective surface density of states of silicon," *Solid-State Electronics*, vol. 17, pp. 335-339, 1974.
- [9] G. Ghibaudo, S. Bruyère, T. Devoivre, B. DeSalvo and E. Vincent, "Improved method for the oxide thickness extraction in MOS structures with ultrathin gate dielectrics," *IEEE Transactions on semiconductor manufacturing*, vol. 13, pp. 152-158, 2000.
- [10] S. Mudanai, F. Li, S. B. Samavedam, P. J. Tobin, C. S. Kang, R. Nieh, J. C. Lee, L. F. Register and S. K. Banerjee, "Interfacial Defect States in HfO₂ and ZrO₂ nMOS Capacitors," *IEEE Electron Device Letters*, vol. 23, no. 12, pp. 728-730, 2002.